

Evaporation-Driven Assembly of Colloidal Particles

Eric Lauga and Michael P. Brenner

Division of Engineering and Applied Sciences, Harvard University, 29 Oxford Street, Cambridge, Massachusetts 02138, USA

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Colloidal particles absorbed at the interface of a liquid droplet arrange into unique packings during slow evaporation [V.N. Manoharan *et al.*, *Science* **301**, 483 (2003)]. We present a numerical and theoretical analysis of the packing selection problem. The selection of a unique packing arises almost entirely from geometrical constraints during the drying.

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Manoharan *et al.* [1] recently presented an ingenious method for fabricating clusters of small particles into precise configurations. Polystyrene spheres (diameter 844 nm) were dispersed in a toluene-water emulsion with each oil droplet containing a low number N of spheres. The toluene was then preferentially evaporated, forcing the particles to come together into compact clusters. Surprisingly, the final particle packings were unique: the observed packings for $N \leq 11$ closely correspond to those previously identified [2] as minimizing the second moment of the particle distribution, $\mathcal{M} = \sum_i \|\mathbf{r}_i - \mathbf{r}_0\|^2$, where \mathbf{r}_0 is the center of mass of the cluster.

The fact that such a simple process leads to precision assembly at the submicron scale points to exciting possibilities for controlling the assembly of more general objects [3]. The goal of this Letter is to understand the physical principles underlying the observations of Manoharan *et al.* [1]. Why are the final packings unique? (Why) do they minimize the second moment? What physical parameters do these results depend upon? We first present numerical simulations of hard spheres on an evaporating liquid droplet for a wide range of liquid-solid contact angles: for each contact angle the simulations reproduce the final packings of [1]. We then demonstrate that the uniqueness of the packings can be understood almost entirely from purely geometrical considerations. The arguments suggest a method for creating new packings, which we confirm through simulations.

Numerical simulations.—The experiments suggest the following theoretical problem: for a given liquid volume, the particle configuration is determined from minimizing the total surface energy

$$U_\Sigma = \gamma_D \int_D dS + \gamma_{DP} \int_{DP} dS + \gamma_P \int_P dS, \quad (1)$$

while respecting excluded volume constraints between the particles. Here D , DP , and P refer to the droplet surface, droplet-particle interface, and particle surface, respectively (see Fig. 1).

Numerical simulations are performed using Brakke's Surface Evolver [4], a program which determines the equilibrium configuration of deformable surfaces given the definition of an energy. The colloidal spheres are

modeled as liquid droplets with high surface tension, typically 1 to 2 orders of magnitude larger than the main droplet, in order to penalize nonspherical deformations of their shape. Interfacial tensions between the droplet and the particles are chosen appropriately in order to satisfy Young's law at the liquid-solid contact line, $\gamma_P = \gamma_D \cos\theta + \gamma_{DP}$, where θ is the equilibrium contact angle. Non-interpenetrability is enforced with an excluded volume repulsion energy U_R acting between the centers of the spheres [5]; U_R dominates when at least two spheres overlap by 1%.

The particles are initially positioned randomly on the droplet. The droplet volume is then slightly decreased (by 1% or less) and the particles rearrange to a new equilibrium, minimum of $\mathcal{U} = U_\Sigma + U_R$. This procedure replaces the evaporation dynamics by a series of equilibrium problems and therefore mimics the low evaporation rate limit of the experiments.

We find that the packings obtained numerically are unique and agree with those obtained by Manoharan *et al.* [1], over the range of contact angles tested ($10^\circ \leq \theta \leq 170^\circ$) and initial conditions. A comparison of the final computational and experimental packings is illustrated in Fig. 2. In all cases, the final second moment obtained numerically differs by less than 0.5% from the minimum moment packings of [2]. The agreement verifies that the experimental packing problem is consistent with the energy minimization posed above and that other particle-particle interactions (electrostatic, van der Waals...) are not essential to generate the final packings.

There is one noteworthy difference between the simulations and the experiments. In the simulations, the clus-

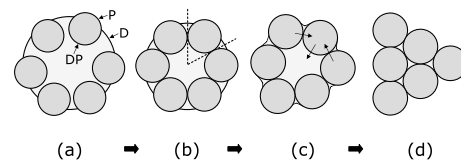


FIG. 1. Schematic two-dimensional representation of the drying experiment: (a) particle configuration when the droplet volume is above the critical volume; (b) critical packing; (c) rearrangement below the critical packing with both capillary and contact forces acting on each particle; (d) final packing.

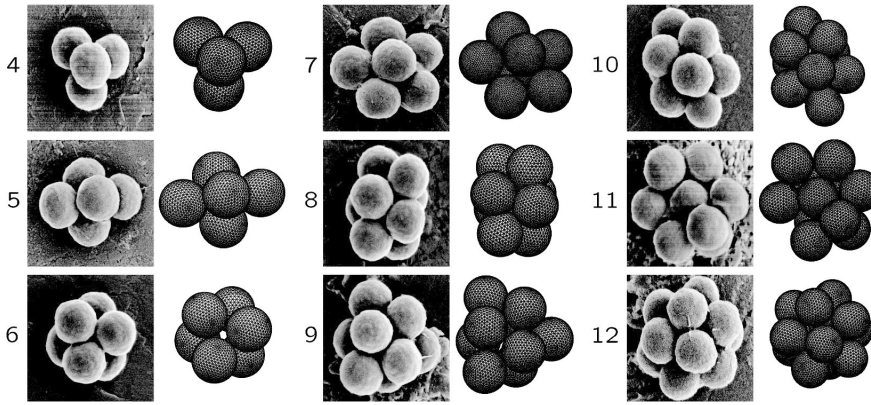


FIG. 2. Comparison of the experimentally observed packings (left) with those obtained by numerical simulations (right) as a function of the number of spherical particles.

ter of spheres evolves smoothly (c.f. Fig. 3). In contrast, experiments [1] showed a discontinuous transition to the final packing. The discrepancies must arise from experimental features not present in the simulations, probably contact angle hysteresis.

The critical volume.—We now study theoretically the packing selection problem. For sufficiently large liquid volumes, the minimum energy solution is a spherical droplet with noninteracting force-free particles [Fig. 1(a)]. There exists a critical volume V_c below which the droplet cannot remain spherical [Fig. 1(b)]; below this critical volume each particle is acted upon by capillary forces. At the critical volume V_c , there is a critical packing of particles, which can be characterized as follows.

The interactions between particles on the surface of a sphere are equivalent to the steric interactions between the “cone of influence” of each particle, defined as the cone originating from the droplet center and tangent to the particle [Fig. 1(b)]. These are also equivalent to the interactions between the intersection of these cones with

the droplet, which are circles. Consequently, packing spherical particles at the critical volume is equivalent to packing circles on a sphere, a mathematical problem with rich history [6–9] and for which numerical solutions have been proposed up to $N = \mathcal{O}(100)$. Usually, the circle packings for a given N are unique, with two types of exceptions. For $N = 5, 19, 20, 23, 26, 28, 29, \dots$ the circle packing has continuous degrees of freedom (where at least one circle is free to “rattle”) [9]; for $N = 15$ two different configurations (15a and 15b) lead to the densest packings with equal surface density [9].

Hence, provided that there are no kinetic traps, identical particles will arrange themselves into the circle packing at the critical volume. Figure 3 confirms this conclusion in numerical simulations for $N = 9$. Changing the contact angle of a particle changes the size of its circle of influence and the droplet radius by the same amount, so the circle packing problem is unchanged and the critical packing of particles at the critical volume is independent of wetting characteristics; the critical droplet volume is, however, contact angle dependent.

How do the particles rearrange when $V < V_c$?—The energy minimization problem suggests that we must find the particle configuration which minimizes U_Σ under interpenetrability and contact angle constraints. Owing to the complexity of solving for the liquid surface of constant mean curvature, this at first appears to be an extraordinarily difficult theoretical problem. However, we have found that, at least at low N , the constraints associated with packing of particles are sufficient to uniquely determine their initial rearrangements.

Let us suppose the droplet volume is reduced by a small amount $\delta V \ll V_c$. Deviations of the droplet interface from spherical lead to capillary forces (\mathbf{F}_i) on each particle. Since every particle must be in force equilibrium, these forces must be balanced by contact forces (\mathbf{f}_{ji}) between the particles [Fig. 1(c)], such that

$$\mathbf{F}_i + \sum_{j \in C(i)} \mathbf{f}_{ji} = \mathbf{0}. \quad (2)$$

where $C(i)$ denotes the set of particles in contact with sphere i [10].

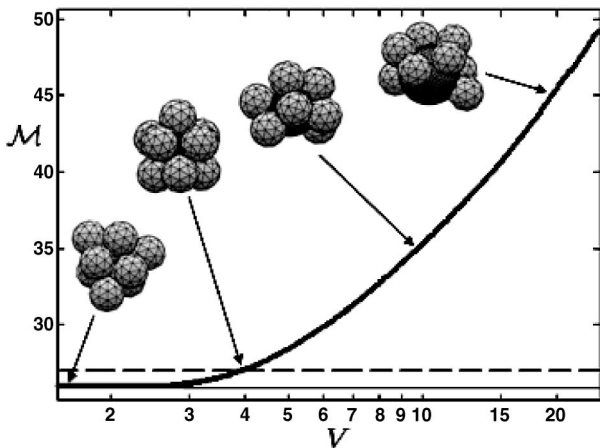


FIG. 3. Numerical evolution of the second moment as a function of the droplet volume for $N = 9$ (both in units of the particle radius). The final second moments are 25.899 (theory) and 25.946 (simulations). Inset, packings observed during the drying process when $\mathcal{M} = 25.946, 27.016, 35.014, 45.058$. Solid line, minimal second moment; dashed line, second moment at the critical packing ($\mathcal{M} = 27$).

Let us now characterize the number of ways the particles can rearrange to accommodate this change in volume. Each particle has 3 degrees of freedom and the droplet has one (the value of its mean curvature or pressure), so there are $3N + 1$ degrees of freedom. The constraints include (a) solid body rotation does not modify the packing (3 constraints); (b) the particles cannot overlap, (n_c constraints, where n_c is the number of contacts at the critical packing) [11]; (c) forces have to balance [Eq. (2)].

Equation (2) implies nontrivial constraints. Suppose we solve (2) for the n_c contact forces $\mathbf{f}_{ji} = f_{ji}\mathbf{e}_{ji}$ ($\mathbf{e}_{ji} = \mathbf{e}_i - \mathbf{e}_j$, where \mathbf{e}_i is the unit vector directed from the droplet center to the particle center). Equation (2) has $3N$ components, and $n_c \sim 2N$ unknowns (Table I). Consequently, solutions exist only if compatibility relations are satisfied between the \mathbf{F}_i . Since capillary forces depend on the position of the particles, these equilibrium considerations constrain the rearrangement of the particles. Close to the critical volume, the capillary forces are given by $\mathbf{F}_i = F_i\mathbf{e}_i$. From the N scalar forces F_i , Eq. (2) shows that only a subset n_f can be chosen independently, *i.e.* equilibrium of each particle leads to $N - n_f$ additional constraints; n_f is found by computing the rank of the compatibility matrix in (2) [12].

The total number of admissible modes of rearrangement n_m for the colloidal particles is found by subtracting the number of constraints from the number of degrees of freedom: we find $n_m = 2N - 2 + n_f - n_c$. For a given N , n_m is entirely determined by the geometry of the critical packing. The results are displayed in Table I. Whenever the circle packing is unique, we find that $n_m = 1$. When $n_m > 1$, we find that the number of modes is always correlated with the presence of rattlers. If a total number of continuous degrees of freedom n_d exist in the circle packing ($n_d = 2$ or one per rattler depending on if it is completely free, as in $N = 19$ or constrained in a slot, as in $N = 5$), we always find that $n_m = 1 + n_d$. Since $n_d \neq 0$ indicates that there exist n_d force-free surface modes for

TABLE I. Characteristics of critical packings of spherical colloidal particles as a function of their number N : number of contacts (n_c), number of independent forces (n_f) and number of admissible modes of rearrangement (n_m).

N	n_c	n_f	n_m	N	n_c	n_f	n_m	N	n_c	n_f	n_m
4	6	1	1	13	24	1	1	21	40	1	1
5	6	1	3	14	28	3	1	22	42	1	1
6	12	3	1	15a	30	3	1	23	43	2	3
7	12	1	1	15b	30	3	1	24	60	15	1
8	16	3	1	16	32	3	1	25	48	1	1
9	18	3	1	17	34	3	1	26	46	1	5
10	19	2	1	18	34	1	1	27	52	1	1
11	25	6	1	19	34	1	3	28	52	1	3
12	30	9	1	20	39	6	5	29	54	1	3

the packing, we obtain therefore that there exists a unique mode of rearrangement for all particles which are non-rattlers. In experiments, the degeneracy in the circle packing problem is chosen by additional information: for example, in [1], because the particles are charged only on the side exposed to the water, there is a weak dipolar repulsive force between the particles which breaks the degeneracy [13,14].

This result implies that there is only a single set of $\{F_i\}$ that is consistent with all the constraints. In five cases ($N = 4, 6, 8, 12, 24$), a decrease of the droplet volume leads to an increase in the droplet pressure without rearrangement of the particles; in these cases, we predict the final packings to be the same as the spherical packings, which is confirmed experimentally and by our simulations for $N = 4, 6, 12$. In all other cases, the modes include rearrangements of the particles and are unique, independent of surface energies and only a function of geometry. Capillarity enters the problem in relating the force F_i to the displacement δr_i . If the radius of curvature of the droplet changes from R to $R + \delta R$ then volume conservation implies that δr_i and δR are related through $A\sum_i \delta r_i + (4\pi R^2 - NA)\delta R = 0$, where A is the wetted area of the particles. The capillary force F_i is then given by

$$F_i = -2\pi\gamma_D \cos\beta \left(\delta r_i + \frac{A \cos\beta}{4\pi R^2 - NA} \sum_i \delta r_i \right), \quad (3)$$

where α is the dry angle of the particle on the critical packing, θ the equilibrium contact angle, and $\beta = \alpha - \theta$. This formula is asymptotically valid in the limit $\delta V \rightarrow 0$ so that deviations from a spherical cap droplet are small.

The above results apply just below V_c . However, the general principle can be iteratively applied below the critical volume: starting from the spherical packing, we decrease the droplet volume by small increments and, assuming Eq. (3) continues to hold as a modeling assumption, we calculate the corresponding incremental particle rearrangement consistent with all the constraints. At each step in the iteration it is necessary to recompute the mode $\{F_i\}$ (and to check whether $f_{ij} > 0$); for every packing there is a unique choice that is consistent with the constraints. This process iterates until the final equilib-

TABLE II. Final second moment of the model (\mathcal{M}_m) compared with the final second moment in the experiment [1] if the particles are assumed to be perfectly spherical (\mathcal{M}_{exp}) and with the minimum second moment (\mathcal{M}_2).

N	\mathcal{M}_{exp}	\mathcal{M}_m	\mathcal{M}_2	N	\mathcal{M}_{exp}	\mathcal{M}_m	\mathcal{M}_2
3	4	4	4	9	25.899	25.899	25.899
4	6	6	6	10	31.828	31.828	31.828
5	9.333	9.333	9.333	11	37.835	37.929	37.835
6	12	12	12	12	43.416	43.416	42.816
7	16.683	17.100	16.683	13	51.316	52.032	47.701
8	21.157	21.657	21.157	14	59.225	61.000	54.878

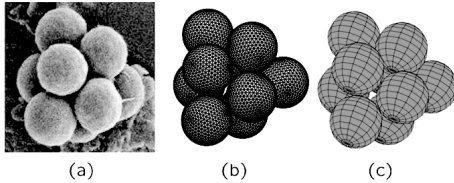


FIG. 4. Comparison for $N = 9$ between the packing given by (a) experiments, (b) simulations, and (c) model.

rium configuration is reached. During the process the packing changes substantially from the initial disk packing—typically multiple new contacts are added.

The results of the model are displayed in Table II and in Fig. 4. The model reproduces accurately the final experimental packings for $N \leq 6$ and $9 \leq N \leq 14$; in particular, the nonconvex packing for $N = 11$ is well predicted by the model. Differences in the final packings for $N = 7$ and 8 likely arise from deviations of the capillary force-particle displacement relationship from the linear law (3) when the change in droplet volume is sufficiently large. It should be emphasized that the computational cost of the new algorithm is orders of magnitude lower than that of a full simulation. To our surprise the final packing can be computed quite accurately without ever knowing the shape of the liquid surface during the packing process!

Finally, we remark on the minimal moment criterion itself: our results suggest that the drying influences the final packings only (a) through enforcing the initial disk packing at the critical volume; and (b) through Eq. (3), relating the F_i to the particle displacements. It is probably not coincidental that (up to prefactors) Eq. (3) is similar to the force-displacement relation $F_i = -\partial_i \mathcal{M} \sim c_1 \delta r_i + c_2 \sum_j \delta r_j$, for the minimal moment criterion. We believe that this explains the similarity between the observed structures and those minimizing the second moment.

The calculations presented herein suggest that the unique packings observed by Manoharan *et al.* arise because (i) the initial circular packing is unique for the regime they explored, except for $N = 15$, and (ii) the subsequent evolution of the particles is so highly constrained that there is only one final packing that is consistent with the constraints.

This suggests that the only way to generate different packings is to modify the circle packing at the critical volume. This can be easily modified by choosing particles with differing sizes or wettabilities (thus creating circles of different sizes). We have run simulations where particles on a given droplet possess different contact angles. Figure 5 shows three cases of $N = 9$ particles with contact angles of either 160° or 20° . The three packings and their second moment differ significantly from Fig. 4.

In summary, we have presented a numerical and theoretical study of the packing selection problem of Manoharan *et al.* [1]. The selection of a unique packing

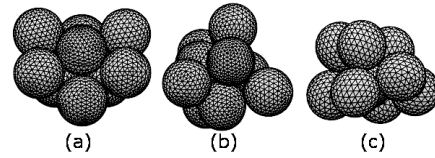


FIG. 5. Final configuration of nine spheres with different wetting conditions: on the droplet, n_1 spheres have a contact angle of 160° and $n_2 = 9 - n_1$ spheres have a contact angle of 20° . (a) cases $\{n_1, n_2\} = \{2, 7\}$ and $\{4, 5\}$, $\mathcal{M} = 27.706$; (b) case $\{n_1, n_2\} = \{6, 3\}$, $\mathcal{M} = 29.780$; (c) case $\{n_1, n_2\} = \{8, 1\}$, $\mathcal{M} = 30.754$ (units of the particle radius).

was found to arise almost entirely from geometrical constraints during the drying process. The analysis suggests that changes in wettability among the particles lead to changes in the final packings.

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- [1] V. N. Manoharan, M. T. Elsesser, and D. J. Pine, *Science* **301**, 483 (2003).
 - [2] N. J. A. Sloane, R. H. Hardin, T. D. S. Duff, and J. H. Conway, *Discrete Comput. Geom.* **14**, 237 (1995).
 - [3] P. A. Kralchevsky and N. D. Denkov, *Curr. Opin. Colloid Interface Sci.* **6**, 383 (2001).
 - [4] K. A. Brakke, *Exp. Math.* **1**, 141 (1992).
 - [5] We have used potentials ranging from exponential $U_R \sim a \exp(1 - d/d_0)$, to power law, $U_R \sim a(d_0/d)^p$, where d (d_0) is the distance between two particles (diameter of a particle). The results are insensitive to the exact functional form of U_R .
 - [6] T. W. Melnyk, O. Knop, and W. R. Smith, *Can. J. Chem.* **55**, 1745 (1977).
 - [7] B. W. Clare and D. L. Kepert, *Proc. R. Soc. London A* **405**, 329 (1986).
 - [8] N. J. A. Sloane, R. H. Hardin, W. D. Smith *et al.*, *Tables of Spherical Codes*, <http://www.research.att.com/njas/packings/>.
 - [9] D. A. Kottwitz, *Acta Crystallogr. Sect. A* **47**, 158 (1991).
 - [10] Torque balance also has to be satisfied; since contact forces do not contribute to any torque, this implies that the capillary force must be directed towards the center of the particles.
 - [11] Here we assume that spheres in contact at the critical packing remain in contact, as suggested by numerical simulations; the consistency of this assumption is verified *a posteriori* by checking that $f_{ij} > 0$.
 - [12] Since $\sum_i \mathbf{F}_i = \mathbf{0}$, an upper bound is $n_f \leq N - 3$.
 - [13] P. Pieranski, *Phys. Rev. Lett.* **45**, 569 (1980).
 - [14] A. J. Hurd, *J. Phys. A* **18**, L1055 (1985).